

Compactifications on Hyperbolic Spaces

Domenico Orlando

CPHT - École polytechnique - Palaiseau (France)
VUB - Bruxelles (Belgium)

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Collaboration with: C.Bachas, C.Kounnas (ENS),
M.Petropoulos (CPHT).



Outline

- 1 Why are we here?
- 2 Preliminaries
- 3 The “AdS Splitting Series”
- 4 Stability
- 5 Source Brane Interpretation
- 6 Conclusions



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Why are we here

Compactifications on Maximally Symmetric Spaces

- Compactifications on maximally symmetric spaces
- Very nice **geometric and algebraic properties**
- Negative curvature **hyperbolic spaces** [Kehagias, Russo]
- Stability of non-supersymmetric backgrounds
- Compactification
- Dual theories
- Cosmology [Townsend]



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Avant-première

Main results

- Hyperbolic space solutions in M-theory
- AdS-splitting
- Hyperbolic solution are stable with respect to small perturbations even without supersymmetry
- M5-brane wrapped around an H_3 manifold



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Short introduction to Hyperbolic Manifolds

Maximally Symmetric Spaces

- A n -dimensional maximally symmetric space can be defined as a **pseudosphere** in $n + 1$ dimensions:

$$\epsilon_0(X^0)^2 + (X^1)^2 + \cdots + (X^{n-1})^2 + \epsilon_n(X^n)^2 = \epsilon L^2$$

$(\epsilon_0, \epsilon_n, \epsilon)$	---	-++	-+-	+++
Space	AdS _n	dS _n	H_n	S^n

Hyperbolic (Poincaré) Spaces

- $H_n = SO(1, n) / SO(n)$ coset.



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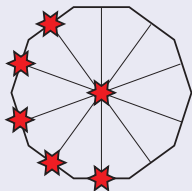
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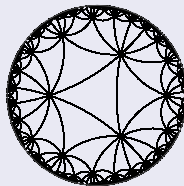


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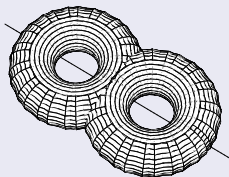
Riemann Surfaces: the Double Torus



Fundamental polygon



Tessellation of the Poincaré disk



Short Introduction to Hyperbolic Manifolds

H_n/Γ : no moduli

- Any finite closed manifold of constant negative curvature is H_n/Γ , $\Gamma \subset SO(1, n)$
- **Rigidity theorem**: the geometry of a finite manifold H_n/Γ is determined by its fundamental group [Mostow]
 - (Algebraic) Given Γ_1 and Γ_2 , lattices in $SO(1, n)$ such as H_n/Γ_i is finite-volume, then if they are isomorphic then they are conjugate.



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H_3/Γ

- Generalization of Riemann surfaces.
- Only lower bounds on the volume $V > 0.166$



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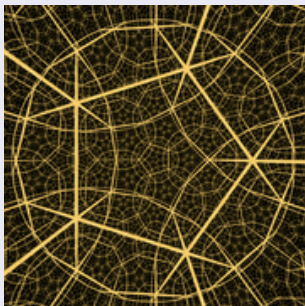


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Example: Seifert-Weber Manifold

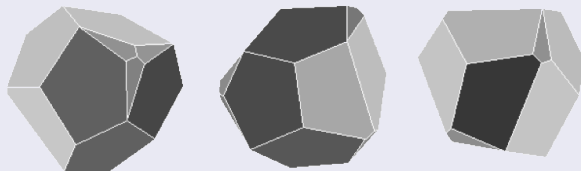


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Example: Weeks Manifold



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M-theory ansatz

Action and Notation

- Ansatz: direct products of symmetric spaces
 $M_{11} = M_0 \times M_1 \times M_2 \times \dots$
- Split the Ricci tensor in blocks

$$R^\mu{}_\nu|_i = k_i \delta^\mu{}_\nu|_i$$

- Choose the gauge fields as:

$$F_I = Q_I \omega_I$$

where $\omega_I = \bigwedge_{i \in I} \omega_i$.



Equations of motion

Algebraic System

- The equation of motion are an **algebraic system** for k_i :

$$2k_i - R = - \sum_I \varepsilon_I(i) \varepsilon_I(0) Q_I^2, \text{ for } i = 0, 1, 2, \dots,$$

where $R = \sum_i d_i k_i$ is the total Ricci scalar and

$$\varepsilon_I(i) = \begin{cases} +1 & \text{if } i \in I, \\ -1 & \text{otherwise.} \end{cases}$$

- Negative values for k_i** are possible due to the stress-energy tensor coming from fields living in other subspaces.
- Only negative contribution on M_0 : **no-go** for de Sitter.



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Splitting Series Solution

AdS Splitting

- We consider **Cartesian products**
- The Ricci tensor has a **block matrix** form
- If each block is a symmetric space there is a **fixed choice** for the curvatures where **AdS₇** is not distinguishable from **AdS_d × H_{7-d}**

$$\text{AdS}_7 \rightarrow \text{AdS}_d \times H_{7-d}$$

- Starting from the well known $\text{AdS}_7 \times S^4$ one gets:

$$\text{AdS}_4 \times H_3 \times S^4$$



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Hyperbolic Compact Manifolds and Supergravity

Killing Spinors on H_n/Γ

- If n is even there is one Killing spinor in $SO(1, n - 1)$
 - To preserve it, $\Gamma \subset SO(1, n - 3)$.
 - H_n/Γ would remain infinite
 - No supersymmetric compactifications for n even
- If n is odd there are two Killing spinors in $SO(1, n - 1)$



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Stability

General Approach

- No protection from supersymmetry. We can check at least the **stability against scalar modes**.
- Turn on **dilatation** of the compact subspaces
- Write the effective, dimensionally-reduced action
- Check the masses for the scalar fields against the **Breitenlohner-Freedman bound**
- In the AdS-splitting case there are **no tachyons**.



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Stability

d -dimensional Action

- Effective action

$$S \sim \int d^d x \sqrt{-g^{(d)}} \left[R^{(d)} - \frac{1}{2} \partial_\mu \Phi_i \partial^\mu \Phi_i - \bar{V}(\Phi) \right]$$

- Φ_1 is the overall volume
- Φ_2 is the ratio between the two volumes
- The effective potential is

$$\bar{V}(\Phi) = e^{-2(d_1\varphi_1 + d_2\varphi_2)/(d-2)} \times \\ \times \left(-e^{-2\varphi_1(x)} R^{(1)} - e^{-2\varphi_2(x)} R^{(2)} + V(\varphi_1, \varphi_2) \right)$$



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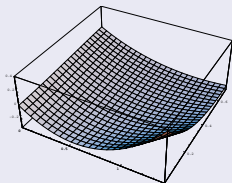
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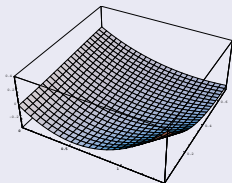
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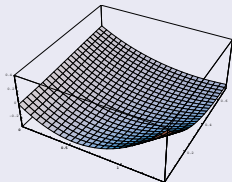
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- Alternative description in terms of **M-branes wrapping the hyperbolic space**
- The stability of the $\text{AdS} \times H \times S$ backgrounds results from the presence of **two competing effects**:
 - The tension given by the M-brane trying to shrink the internal manifold
 - The repulsion due to the negative cosmological constant of the hyperbolic space
- Only one magnetic charge. Back-reaction of an M5-brane



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M5 brane

M5-brane Ansatz

- M5 brane wrapping a H_3 hyperbolic manifold
- Most simple ansatz, preserving $SO(1,2) \times SO(4) \times SO(1,3)$

$$ds^2 = e^{2\varphi_1(r)} \left(-dx_0^2 + dx_1^2 + dx_2^2 \right) + e^{2\varphi_2(r)} \left(dr^2 + r^2 d\Omega_4^2 \right) + e^{2\varphi_3(r)} \left(\frac{dx_8^2 + dx_9^2 + dx_{10}^2}{x_8^2} \right)$$



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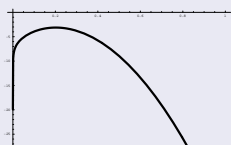
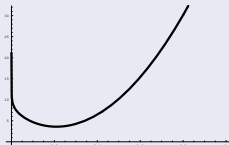
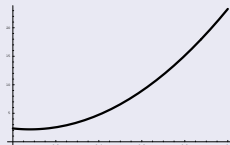
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Numerical Solution

- The system can be solved numerically


 $\varphi_1(r)$

 $\varphi_2(r)$

 $\varphi_3(r)$

- Small- r regime: $\text{AdS}_4 \times S^4 \times H_3$ solution

$$\varphi_1(r) \sim \log r \quad \varphi_2(r) \sim -\log r \quad \varphi_3(r) = \text{const.}$$

- Asymptotically flat

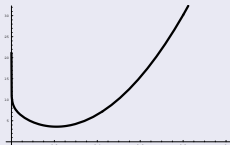
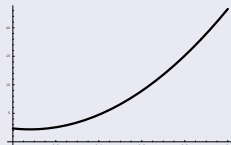
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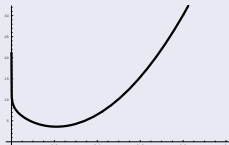
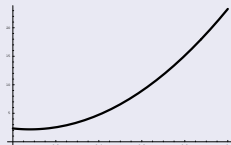
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Outline

- 1 Why are we here?
- 2 Preliminaries
- 3 The “AdS Splitting Series”
- 4 Stability
- 5 Source Brane Interpretation
- 6 Conclusions**



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- Hyperbolic solution stable with respect to small perturbations even without supersymmetry
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to boldly go...



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